

Two-Loop Corrections to Higgs Boson Masses in the Complex MSSM

Sebastian Paßehr

in collaboration with Prof. Wolfgang Hollik

Max Planck Institute for Physics, Munich

Trieste,

29th of August 2013



MAX-PLANCK-GESELLSCHAFT



- ① Higgs Bosons
- ② Higher Order Corrections
- ③ Order α_t^2 Corrections
- ④ Status and Outlook

Higgs bosons in the MSSM

two complex $SU(2)$ -Higgs doublets necessary,

$$h_1 = \begin{pmatrix} \frac{1}{\sqrt{2}} (v_1 + \phi_1^0 - i\gamma_1^0) \\ -\phi_1^- \end{pmatrix} \quad \text{and} \quad h_2 = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (v_2 + \phi_2^0 - i\gamma_2^0) \end{pmatrix},$$

\Rightarrow eight bosonic degrees of freedom:

3 Goldstone bosons, 5 physical Higgs bosons,

real parameters:

Goldstone bosons	$G^0, G^\pm,$
physical CP even bosons	$h^0, H^0,$
physical CP odd bosons	$A^0,$
physical charged Higgs bosons	$H^\pm.$

- Higgs potential fixed by Lagrangian density:

$$V_{\text{Higgs}} = m_1^2 h_1^\dagger h_1 + m_2^2 h_2^\dagger h_2 - m_{12}^2 (h_1 \cdot h_2 + h_1^\dagger \cdot h_2^\dagger) \\ + \frac{1}{8} (g_1^2 + g_2^2) (h_2^\dagger h_2 - h_1^\dagger h_1)^2 + \frac{1}{2} g_2^2 h_1^\dagger h_1 h_2^\dagger h_2,$$

- tree-level masses correlated:

$$m_{H^0, h^0}^2 = \frac{1}{2} \left(m_{A^0}^2 + m_Z^2 \pm \sqrt{(m_{A^0}^2 + m_Z^2)^2 - (2m_Z m_{A^0} \cos 2\beta)^2} \right), \\ m_{H^\pm}^2 = m_{A^0}^2 + m_W^2,$$

- two free parameters: conventionally $\tan \beta = \frac{v_2}{v_1}$, m_{A^0} ,
- theoretical upper bound: $m_{h^0}^2 \leq (m_Z \cos 2\beta)^2$.

Higgs boson masses in higher orders

real parameters:

$$\mathcal{M} = \begin{pmatrix} m_{h^0}^2 - \hat{\Sigma}_{H^0 H^0} & \\ & m_{H^0}^2 - \hat{\Sigma}_{H^0 H^0} \end{pmatrix}, \quad \hat{\Sigma}_{A^0 A^0} := 0,$$

Higgs boson masses in higher orders

real parameters:

$$\mathcal{M} = \begin{pmatrix} m_{h^0}^2 - \hat{\Sigma}_{h^0 h^0} & -\hat{\Sigma}_{h^0 H^0} \\ -\hat{\Sigma}_{H^0 h^0} & m_{H^0}^2 - \hat{\Sigma}_{H^0 H^0} \end{pmatrix}, \quad \hat{\Sigma}_{A^0 A^0} := 0,$$

Higgs boson masses in higher orders

complex parameters ($\phi_{A_t}, \phi_\mu, \dots$):

$$\mathcal{M} = \begin{pmatrix} m_{h^0}^2 - \hat{\Sigma}_{h^0 h^0} & -\hat{\Sigma}_{h^0 H^0} & -\hat{\Sigma}_{h^0 A^0} \\ -\hat{\Sigma}_{H^0 h^0} & m_{H^0}^2 - \hat{\Sigma}_{H^0 H^0} & -\hat{\Sigma}_{H^0 A^0} \\ -\hat{\Sigma}_{A^0 h^0} & -\hat{\Sigma}_{A^0 H^0} & m_{A^0}^2 - \hat{\Sigma}_{A^0 A^0} \end{pmatrix}, \quad \hat{\Sigma}_{H^\pm H^\pm} := 0,$$

(most general case also includes longitudinal G^0 and Z),

Higgs boson masses in higher orders

complex parameters ($\phi_{A_t}, \phi_\mu, \dots$):

$$\mathcal{M} = \begin{pmatrix} m_{h^0}^2 - \hat{\Sigma}_{h^0 h^0} & -\hat{\Sigma}_{h^0 H^0} & -\hat{\Sigma}_{h^0 A^0} \\ -\hat{\Sigma}_{H^0 h^0} & m_{H^0}^2 - \hat{\Sigma}_{H^0 H^0} & -\hat{\Sigma}_{H^0 A^0} \\ -\hat{\Sigma}_{A^0 h^0} & -\hat{\Sigma}_{A^0 H^0} & m_{A^0}^2 - \hat{\Sigma}_{A^0 A^0} \end{pmatrix}, \quad \hat{\Sigma}_{H^\pm H^\pm} := 0,$$

(most general case also includes longitudinal G^0 and Z),

meaning of $\hat{\Sigma}$:

$$\hat{\Sigma} = \Sigma(p^2) - \delta m^2 = \hat{\Sigma}(p^2),$$

$$\hat{\Sigma}(p^2) = \hat{\Sigma}^{\text{one loop}}(p^2) + \hat{\Sigma}^{\text{two loop}}(p^2) + \dots$$

- main contributions come from t and \tilde{t} loops; order α_t , but proportional to m_t^4 :

$$\Sigma_{hh} = \text{---} h^0 \text{---} \text{---} \text{---} h^0 \text{---} + \text{---} h^0 \text{---} \text{---} \text{---} h^0 \text{---} + \text{---} h^0 \text{---} \text{---} \text{---} h^0 \text{---},$$

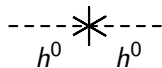
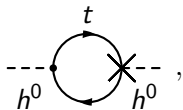
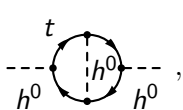
- additional parameters: $m_{\tilde{q}_L}, m_{\tilde{t}_R}, A_t, \mu,$
- mass contribution: ca. 40% of tree-level result,
- uncertainty of the calculation still too big (experimental accuracy $\approx \pm 0.6\text{GeV}$).

most important parts:

corrections to m_t -enhanced one-loop contributions
in a gauge-less limit,

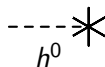
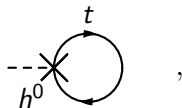
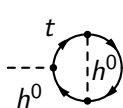
- corrections by gluons and gluinos already known,
order $\alpha_t \alpha_s$ in an on-shell scheme,
[Heinemeyer, Hollik, Rzehak, Weiglein, arXiv:hep-ph/0705.0746, 2007],
- corrections by Higgs and Higgsinos already known
in the real MSSM in the effective potential approach,
order α_t^2 in a $\overline{\text{DR}}$ scheme,
[Brignole, Degrassi, Slavich, Zwirner, arXiv:hep-ph/0112177, 2002],
- corrections by Higgs and Higgsinos
in the case of the complex MSSM: in process.

- again: enhancement by additional m_t^2 ,
- Feynman-diagrammatic approach:



(two-loop) (one-loop) $\cdot (\delta^{(1)})$, $(\delta^{(2)}) + (\delta^{(1)})^2$,
 $+ (\text{one-loop})^2$,

- $(\delta^{(2)}) + (\delta^{(1)})^2$ acquired by renormalizing the Higgs potential, calculation of tadpole diagrams necessary:



- creation of Feynman-diagrams and amplitudes with FeynArts,
[Hahn, arXiv:hep-ph/0012260, 2001],
- applying approximations,
- reducing one-loop diagrams to master integrals with FormCalc,
[Hahn, arXiv:hep-ph/0901.1528, 2009],
- reducing two-loop diagrams to master integrals with TwoCalc,
[Weiglein, Scharf, Böhm, arXiv:hep-ph/9310358, 1993],
- creating counterterms from the Higgs potential,
- applying renormalisation scheme,
- evaluating renormalisation constants with FeynArts and FormCalc,
- expanding master integrals, simplifying result.

(similar as for $\alpha_t \alpha_s$ corrections)

- ① gauge-less limit: $g_1 = 0, g_2 = 0$,
 - only Yukawa-couplings left,
 - $m_W = 0, m_Z = 0$,
 - $m_{h^0} = 0, m_{G^0} = 0, m_{G^\pm} = 0, m_{H^0} = m_{H^\pm}, m_{A^0} = m_{H^\pm}$,
 - $m_{\tilde{\chi}_{3,4}^0} = \mu, m_{\tilde{\chi}_2^\pm} = \mu$,
 - other Charginos and Neutralinos decouple,
 - Higgs mixing angle $\alpha = \beta - \frac{\pi}{2} < 0$,
- ② external momentum equal to zero,
 - only two-loop vacuum diagrams; known analytically,
 - renormalisation constants for genuine two-loop counterterms calculated at zero momentum,
- ③ bottom mass equal to zero,
 - no mixing in sbottom sector,
 - one sbottom (w. l. o. g. \tilde{b}_2) decouples,
 - $m_{\tilde{b}_1}^2 = m_{\tilde{t}_1}^2 - m_t^2$.

required renormalisation constants:

① at one-loop:

- δm_t , $\delta m_{\tilde{t}_1}$ and $\delta m_{\tilde{t}_2}$ fixed by on-shell condition,
- δm_{b_1} dependent on top-stop-sector,
- δA_t fixed by on-shell condition for mixing of stops,
- $\delta \mu$ fixed by on-shell condition for $\tilde{\chi}_2^\pm$ or defined in $\overline{\text{DR}}$ scheme,
- Higgs field renormalisation constants
 $\delta Z_{H_1}|_{\text{div.}}$ and $\delta Z_{H_2}|_{\text{div.}}$ defined in $\overline{\text{DR}}$ scheme,
- $\delta \tan \beta = \frac{\tan \beta}{2} (\delta Z_{H_2} - \delta Z_{H_1})|_{\text{div.}}$ defined in $\overline{\text{DR}}$ scheme,
- δM_W , δM_Z (m_t^2 -dependent part of $\frac{\delta M_W}{M_W}$ and $\frac{\delta M_Z}{M_Z}$),

② at one-loop and two-loop:

- tadpoles δt_{h^0} , δt_{H^0} , δt_{A^0} fixed by on-shell conditions,
- δm_{H^\pm} fixed by on-shell condition,
- δm_{h^0} , δm_{H^0} , δm_{A^0} and $\delta m_{h^0 H^0}$, $\delta m_{h^0 A^0}$, $\delta m_{H^0 A^0}$
 dependent on tadpoles, δm_{H^\pm} and $\delta \tan \beta$, e. g.

$$\delta^{(2)} m_{h^0}^2 =$$

$$-\frac{e}{2M_W s_W} \left[\delta^{(2)} t_{h^0} + c_\beta^2 \delta \tan \beta \cdot \delta t_{H^0} - \delta t_{h^0} \left(\frac{\delta M_W}{M_W} + \frac{\delta s_W}{s_W} \right) \right] + \dots$$

current status:

- all Feynman diagrams are generated and calculated,
- renormalisation constants for complex parameters determined,
- all divergencies cancel,
- first numerical checks:
 $\mathcal{O}(1\text{GeV})$ corrections by two-loop diagrams

outlook:

- inclusion into FeynHiggs,
[Hahn, Heinemeyer, Hollik, Rzehak, Weiglein, arXiv:hep-ph/1007.0956, 2010],
- investigate influence of external momentum.