Four dimensional supersymmetric Yang-Mills quantum mechanics with SU(3) gauge group.

Zbigniew Ambroziński

Jagiellonian University, Krakow, Poland Max Planck Institute, Potsdam, Germany

zbigniew.ambrozinski@uj.edu.pl

29.08.2013

work done in collaboration with P. Korcyl and J. Wosiek







The model

The hamiltonian is given by

$$\begin{split} H &= \frac{1}{2} p_a^i p_a^i + \frac{g^2}{4} f_{abc} f_{ade} x_b^i x_c^j x_d^i x_e^j + \frac{ig}{2} f_{abc} \psi_a^T \Gamma^i \psi_b x_c^i \\ &i, j = 1, \dots, D-1 \text{ - spatial indices} \\ &a, b, c, d, e = 1, \dots, N^2 - 1 \text{ - color indices} \\ &\psi_{a,\alpha} \text{ - Majorana spinor} \end{split}$$

H is supersymmetric. Supersymmetry generators are:

$$Q_{\alpha} = (\Gamma^{i}\psi_{a})_{\alpha}p_{a}^{i} + igf_{abc}(\Sigma^{ij}\psi_{a})_{\alpha}x_{b}^{i}x_{c}^{j}, \qquad (1)$$

where Γ^k - alpha matrices, $\Sigma^{ij}=-\frac{i}{4}[\Gamma^i,\Gamma^j].$

$$\{Q_{\alpha}, Q_{\beta}^{\dagger}\} = 4H\delta_{\alpha\beta} \tag{2}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

In our case D = 4, N = 3.

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

Motivations

BFSS conjecture

uncompactified 11 dimensional M-theory \Leftrightarrow large N limit of supersymmetric quantum mechanics in 10 dimensions [BFSS]

small volume approximation to QCD

bosonic sector of considered model is 0 - order approximation of QCD in small volume approach (i.e.: dynamics of homogeneous fields)

Method

Earlier results - overview



- analytical and numerical results for D = 2 and arbitrary N [Trzetrzelewski; Korcyl]
- numerical results for D = 4, $N \le 6$ in bosonic sector only as 0-order approximation to Yang-Mills QFTs [Lüscher; Ziemann]
- numerical results for *D* = 4, *N* = 2 [Wosiek, Campostrini]
- numerical results at finite tempetature [Catterall, Wiseman; Anagnostopoulos et al.]

◆□▶ ◆□▶ ◆□▶ ◆□▶ ●□



- construct Fock space of gauge invariant states with cutoff on number of bosonic excitations
- construct matrices of hamiltonian, angular momentum and supersymmetry generators
- for energies diagonalize the hamiltonian
- fermionic number is conserved (for D = 2, 4) consider each fermionic sector separately

• rotation symmetry - use sectors of definite angular momentum

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

What do we obtain

Model

- energies (with distinction between continuous and discrete spectrum)
- energy eigenstates (with definite fermion number and angular momentum)
- $\bullet \ \ \text{supersymmetric fractions} \rightarrow \text{identifying supermultiplets}$
- restricted Witten index (sum over sectors with *n_F* even)

Method

Results ●000000

Energies in bosonic sector



・ロト ・聞ト ・ヨト ・ヨト

3

Very fast convergence of lowest energies \Rightarrow spectrum is discrete.

 Model
 Method
 Results

 000
 0
 0

 Digression: signature of continuous spectrum

Spectrum is continuous for kinetic energy only: $H = \frac{1}{2}p_a^i p_a^i$.



▲□▶ ▲圖▶ ★ 国▶ ★ 国▶ - 国 - のへで

Energy behavior is $E \sim 1/N_{cut}$.

Method

Results 00●0000

E 990

Energies in sector with one fermion



・ロト ・聞ト ・ヨト ・ヨト

Still only discrete spectrum.

Method

Results 000●000

Energies in sector with two fermions



・ロト ・聞ト ・ヨト ・ヨト

Spectrum is rather discrete.

Method

Results 0000●00

E 990

Energies in sector with three fermions



・ロト ・四ト ・ヨト ・ヨト

Candidate for continuous spectrum.

Method

Results 00000●0

Supersymmetric fractions

$$\begin{aligned} \mathcal{Q}_{1/2}^{\dagger} &= \frac{1}{2} (Q_1 - iQ_2 + Q_3 + iQ_4) \\ \mathcal{Q}_{-1/2}^{\dagger} &= \frac{1}{2} (iQ_1 + Q_2 - iQ_3 + Q_4) \end{aligned} \qquad \qquad H = \frac{1}{4} \{ \mathcal{Q}, \mathcal{Q}^{\dagger} \} \end{aligned}$$

 $|n_F jmE\rangle$ - eigenstate of hamiltonian

$$(2j+1) = \frac{1}{E} \langle n_F j m E | H | n_F j m E \rangle$$

=
$$\sum_{j'E'} \left(\underbrace{\frac{1}{4E} \sum_{mm'} |\langle (n_F - 1)j'm'E' | Q | n_F j m E \rangle|^2}_{\text{supersymmetric fraction } q_{n_F}(j'E' | jE)} \right)$$

supermultiplets form diamonds like that on the right



Method

Results 000000

Overall picture - identification of supermultiplets





Our results vs earlier results of [Campostrini, Wosiek] for SU(2). Single lines mean that the whole supermultiplet was not identified.

SU(2)

SU(3)

◆ロト ◆昼 → ◆ 臣 → ◆ 臣 → のへぐ

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

Conclusions

- our method gives us a qualitative picture of the spectrum
- the cutoff is still too low (for larger number of fermions) to determine continuous spectrum and to calculate Witten index

▲ロト ▲帰 ト ▲ ヨ ト ▲ ヨ ト ・ ヨ ・ の Q ()

Conclusions

- our method gives us a qualitative picture of the spectrum
- the cutoff is still too low (for larger number of fermions) to determine continuous spectrum and to calculate Witten index

Main challenge

size of basis

- $\bullet\,$ at present we have matrices of sizes up to $4k \times 4k$
- getting nb=7 in each fermion sector would require matrices $1M \times 1M$
- possible solution for higher *D* or *N*: take only most significant basis states?

Literature

- T. Banks, W. Fischler, S.H. Shenker, L. Susskind, *M theory as a matrix model: A Conjecture*, Phys.Rev. D55 (1997) 5112-5128
- M. Trzetrzelewski, Large N behavior of two dimensional supersymmetric Yang-Mills quantum mechanics, J.Math.Phys. 48 (2007) 012203;
 P. Korcyl, Solutions of D=2 supersymmetric Yang-Mills quantum mechanics with SU(N) gauge group, J.Math.Phys. 52 (2011) 052105
- M. Campostrini, J. Wosiek, High precision study of the structure of D=4 supersymmetric Yang-Mills quantum mechanics, Nucl.Phys. B703 (2004) 454-498
- M. Lüscher, G. Münster, Weak Coupling Expansion of the Low Lying Energy Values in the SU(2) Gauge Theory on a Torus, Nucl.Phys. B232 (1984) 445;
 V. Ziemann, Qualitative Untersuchung des niedrig liegenden Spektrums reiner Yang-Mills Theorien im endlichen Volumen mit besonderer Bercksichtigung von SU(3), Diploma Thesis, Universitt Hamburg, 1986
- S. Catterall, T. Wiseman Black hole thermodynamics from simulations of lattice Yang-Mills theory, Phys.Rev.D 78 (2008) 041502;
 S. Catterall, T. Wiseman, Towards lattice simulation of the gauge theory duals to black holes and hot strings, JHEP 0712 (2007) 104
- K. N. Anagnostopoulos, M. Hanada, J. Nishimura and S. Takeuchi, *Monte Carlo studies of supersymmetric matrix quantum mechanics with sixteen supercharges at finite temperature*, Phys.Rev.Lett. 100 (2008) 021601